

Bohr Radius

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In today's lecture we proposed a simple way to find the Bohr radius a_0 . The main idea is to consider the energy of the electron in the frame of the inertial frame of the proton

$$E = T + V = \frac{p^2}{2m} - \frac{kq^2}{a_0}$$

and then set $\frac{\partial E}{\partial a_0} = 0$ and solve for a_0 . An intermediate consideration was the position-momentum uncertainty principle, which helped us write $\frac{p^2}{2m} \sim \frac{\hbar^2}{2ma_0^2}$. But I think the application of the uncertainty principle is tenuous, since the uncertainty principle is really saying $\Delta a_0 \Delta p \sim \hbar$, not $a_0 p \sim \hbar$.

The variational principle offers a more rigorous calculation of the Bohr radius:

Using some trial wave function $\psi(r) = Ae^{-\alpha r}$, where $\alpha \equiv \frac{1}{a}$, we can normalize to find A:

$$\begin{aligned} 1 &= \int \int \int |\psi|^2 d\vec{r} = 4\pi \int_0^\infty r^2 |\psi|^2 dr \\ &= 4\pi A^2 \int_0^\infty r^2 e^{-2\alpha r} dr \\ &= \frac{A^2 \pi}{\alpha^3} = 1 \end{aligned}$$

$$\text{so we find } A = \sqrt{\frac{\alpha^3}{\pi}}$$

Now we seek to find some α_{optimal} such that $\langle H \rangle(\alpha_{\text{optimal}})$ is a minimum. Note

$$\langle H \rangle(\alpha) = \langle V \rangle(\alpha) + \langle T \rangle(\alpha)$$

Let's find the expectation values of the potential and kinetic energies:

$$\begin{aligned} \langle V \rangle &= \int \int \int V(r) |\psi|^2 d\vec{r} \\ &= -4\pi \int r^2 \frac{kq^2}{r} e^{-2\alpha r} \end{aligned}$$

$$\text{integration...} = -kq^2\alpha$$

Meanwhile,

$$\begin{aligned}\langle T \rangle &= \int \int \int \psi^* \left(\frac{\hbar}{2m} \frac{\partial^2}{\partial r^2} \right) \psi \, d\vec{r} \\ \text{by parts:} &= \frac{4\pi\hbar^2}{2m} \int r^2 \left| \frac{\partial \psi}{\partial r} \right|^2 \, dr \\ &= \frac{4\pi\hbar^2}{2m} A^2 \int r^2 \frac{\partial e^{-2\alpha r}}{\partial r} \, dr \\ \text{integration...} &= \frac{\hbar\alpha^2}{2m}\end{aligned}$$

So

$$\langle H \rangle(\alpha) = \langle V \rangle(\alpha) + \langle T \rangle(\alpha) = -kq^2\alpha + \frac{\hbar\alpha^2}{2m}$$

We want α_{optimal} such that $\langle H \rangle(\alpha_{\text{optimal}})$ is a minimum:

$$\frac{\partial \langle H \rangle}{\partial \alpha} = -kq^2 + \frac{k^2\alpha}{m} = 0$$

Solving, we find that $\alpha_{\text{optimal}} = \frac{kq^2m}{\hbar^2}$. But $\alpha \equiv \frac{1}{a}$, so the Bohr radius is

$$a_0 = \frac{\hbar}{kq^2m}$$

Incidentally,

$$\langle H \rangle(\alpha_{\text{optimal}}) = \langle V \rangle(\alpha_{\text{optimal}}) + \langle T \rangle(\alpha_{\text{optimal}}) = \frac{k^2q^4m}{2\hbar^2} - \frac{k^2q^4m}{\hbar^2} = -\frac{k^2q^4m}{2\hbar^2} = -R_y$$

where R_y is the Rydberg energy. It is nice that $\frac{1}{2}\langle V \rangle = \langle T \rangle$. This is what we expect from the virial theorem.