## The Evanescent Spectrum of a Spherical Wave

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In Prof. Hamilton's Wave Phenomena class, we showed that the angular spectrum of a spherical wave (i.e., its decomposition into plane waves) is

$$
\frac{e^{ikr}}{r} = \frac{i}{2\pi} \iint_0^\infty \frac{e^{ik_z|z|}}{k_z} e^{i(k_x x + k_y y)} dk_x dk_y, \qquad k_z = \sqrt{k^2 - k_x^2 - k_y^2}.\tag{1}
$$

As the field is axisymmetric about the z-axis, this is equivalent to

$$
\frac{e^{ikr}}{r} = i \int_0^\infty \frac{e^{ik_z|z|}}{k_z} J_0(\kappa \rho) \kappa d\kappa \tag{2}
$$

where

$$
\kappa = \sqrt{k_x^2 + k_y^2}, \qquad \rho = \sqrt{x^2 + y^2}, \qquad k_z = \sqrt{k^2 - \kappa^2}, \qquad r = \sqrt{\rho^2 + z^2}.
$$
\n(3)

We were asked: What is the interpretation of the evanescent spectrum where  $\kappa > k$ ? I may not have a totally satisfying answer, but I have some thoughts...

To gain intuition we examine the field along the positive  $z$  axis, where  $(2)$  reduces to

$$
\frac{e^{ikz}}{z} = i \int_0^\infty \frac{e^{ik_z z}}{k_z} \kappa d\kappa.
$$
 (4)

We split this integral into the propagating and evanescent components:

$$
\frac{e^{ikz}}{z} = i \int_0^k \frac{e^{i\sqrt{k^2 - \kappa^2}z}}{\sqrt{k^2 - \kappa^2}} \kappa d\kappa + \int_k^\infty \frac{e^{-\sqrt{\kappa^2 - k^2}z}}{\sqrt{\kappa^2 - k^2}} \kappa d\kappa.
$$
 (5)

Both of these integrals are evaluated with u-substitution, yielding

$$
\mathcal{I}_P = i \int_0^k \frac{e^{i\sqrt{k^2 - \kappa^2}z}}{\sqrt{k^2 - \kappa^2}} \kappa d\kappa = \frac{e^{ikz} - 1}{z}, \qquad \mathcal{I}_E = \int_k^\infty \frac{e^{-\sqrt{\kappa^2 - k^2}z}}{\sqrt{\kappa^2 - k^2}} \kappa d\kappa = \frac{1}{z}.
$$
 (6)

In the limit  $z \to 0$  the propagating components yield the finite value  $\mathscr{I}_P \to ik$ . It seems the propagating components are unable to reproduce the singularity at  $z = 0$ , and therefore the evanescent components  $\mathscr{I}_E$  are needed.

This is similar to what Brekhovskikh wrote ([1], p. 231):

It is impossible to obtain the field which would have the required singularity as  $R \to 0$ by superposition of ordinary plane waves only... We have waves propagating in the horizontal plane (the xy plane) with a wavelength approaching zero and simultaneously attenuating in the vertical direction with an attenuation coefficient approaching infinity. At  $x = y = 0$  the superposition of an infinite number of these waves gives an infinite value for the field.

## References

[1] L. M. Brekhovskikh, Waves in Layered Media, 2nd ed. (1980).